## Bayes vs the virial theorem

John Magorrian

Gaia Challenge Workship, Surrey, 20 August 2013 "superGaia": Gaia w/o observational errors, selection effects.

#### An idealised problem

superGaia gives us  $(\mathbf{x}, \mathbf{v})$  for every star in the Galaxy. What's the gravitational potential in which they move?

**Assume:** collisionless, 
$$\partial_t = 0 \rightarrow 2$$
 functions + Jeans thm

$$\Phi(\mathbf{x}), \qquad f(\mathbf{x}, \mathbf{v}) = f(\mathbf{J}), \qquad \mathbf{J} = \mathbf{J}(\mathbf{x}, \mathbf{v}|\Phi)$$
Want this  $\mathbf{J}$  (observations are a **discrete realisation** of this

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$$\int \int \left\{ \begin{array}{c} \text{observations are a discrete} \\ \text{realisation of this} \end{array} \right\}$$

### Standard methods don't work (well)

#### I. Parametrized DF

- Guess some form for  $f(\mathbf{J}|\alpha)$ , parametrised by  $\alpha$
- Try a Φ
- Marginalise or fiddle parameters  $\alpha$
- Assign likelihood to Φ.

But what if your guess for  $f(\mathbf{J}|\alpha)$  is wrong? (Ex: real galaxy has streams/substructure...)

#### II. Schwarzschild's method

- Pick a trial potential Φ
- Launch orbit from each observed  $(\mathbf{x}_n^{\star}, \mathbf{v}_n^{\star})$  in  $\Phi$
- ermm..

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CBE is

$$\mathbf{v} \cdot \frac{\partial f}{\partial \mathbf{x}} = \frac{\partial \Phi}{\partial \mathbf{x}} \cdot \frac{\partial f}{\partial \mathbf{v}}.$$

If we knew  $f(\mathbf{x}, \mathbf{v})$  perfectly then moment-based methods (Jeans, virial) would cough up  $\Phi$  very easily.

## But

All we have is *N* stars drawn from *f*:

- We don't know  $f(\mathbf{x}, \mathbf{v})$  itself,
- nor do we know its moments particularly well!
- Want to avoid taking derivatives of moments, or solving awkward implicit relns among moments.

## Another way of tackling the problem

Motivation: Shirley, there must be a Bayesian version of the (generalised) virial theorem...

(unreadable arXiv:1303.6099)

# Broad-brush restatement of problem (JM, arXiv:1303.6099)

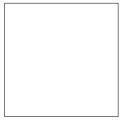
Problem:

- We have two unknown functions,  $\Phi(\mathbf{x})$  and  $f(\mathbf{x}, \mathbf{v})$ .
- Our data *D* is a discrete sample of *f*, but want  $\Phi$ .
- Jeans:  $f(\mathbf{x}, \mathbf{v}) = f(\mathbf{J})$ , where  $\mathbf{J} = \mathbf{J}(\mathbf{x}, \mathbf{v} | \Phi)$ .

Solution: Marginalise *f*!

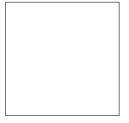
• But how to represent it?

new idea alert!



#### Have: unit stick of starlight + big action-space box.

- Snap off a chunk  $\pi_k$  of starlight
  - Draw  $\beta_k \sim \text{Beta}(1, \alpha)$  and set  $\pi_k = \beta_k \times [\text{stick length}]$
- It is chunk  $\pi_k$  into the action-space box
  - where it lands given by uniform prior in J
- 3 Smear out this  $\pi_k$  in action space:
  - Spike of height  $\pi_k$  at some point  $\mathbf{J}_{\text{land}}$
  - convolve with Gaussian having randomly chosen inverse covariance matrix drawn from uninformative prior (Wishart)
- Apply steps 1–4 to remainder of stick.



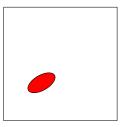
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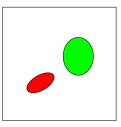
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#### Dirichlet process mixture = distribution of blobs

DF is composed of K blobs ("streams"?), each having

- some probability mass  $\pi_k$ , (k = 1...K)
- centred on some J<sub>k</sub>, and
- with inverse covariance  $\Lambda_k$  (size/shape).

Single free hyperparameter  $\alpha$  controls clumpiness of DF.

Different ways of thinking about Dirichlet process:

- Stick-breaking construction
- Chinese Restuarant process
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#### How to apply it to constraining $\Phi$

Suppose we have a big list of potentials  $\Phi$  to try. Given a trial  $\Phi$ , each  $(\mathbf{x}_n^*, \mathbf{v}_n^*) \rightarrow (\mathbf{J}_n^*, \theta_n^*)$ . Then given a DF  $f(\mathbf{J})$  the likelihood is simply

$$\mathbf{p}(D|\Phi,f)=\prod_{n=1}^N f(\mathbf{J}_n^{\star}).$$

**Easy!** Sum  $p(D|\Phi, f)$  over all *f* generated by the stick-breaking procedure! Obtain **marginalised likelihood**  $p(D|\Phi)$ .

$$p(D|\Phi) = p(\mathbf{J}_{1}^{\star}, ..., \mathbf{J}_{N}^{\star})$$
$$= \sum_{f} p(D|\Phi, f) = \int dp(f) \prod_{n=1}^{N} f(\mathbf{J}_{n}^{\star})$$
$$= \text{ginormous sum over partitions of } N \text{ period}$$

**NB:**  $p(D|\Phi)$  depends only on how the  $J_n^*$  are distributed!

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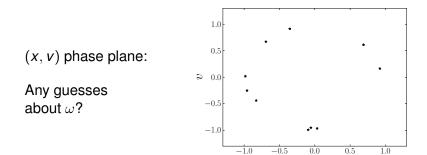
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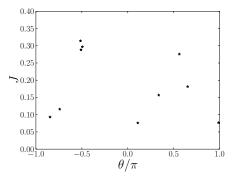
Ptles moving in 1d SHO potential,  $\Phi(x) = \frac{1}{2}\omega^2 x^2$ .



x

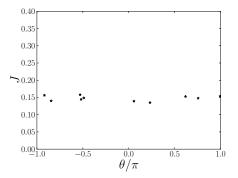
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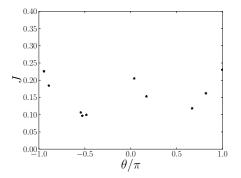
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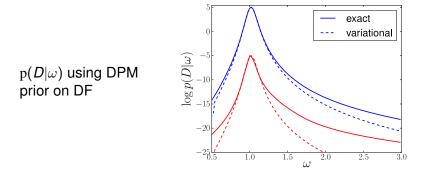


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 $p(D|\Phi)$  is biggest for  $\Phi$  that make J distn "sharpest": - DPM naturally "likes" sharply peaked, clumpy DFs!

Plot shows two different schemes for calculating  $p(D|\Phi)$ .

## Aside: relation to "minimum entropy" idea

(Peñarrubia, Koposov, Walker 2012, also Helmi talk)

A galaxy's entropy,

$$S[f] = -\int f \log f \,\mathrm{d}^3 \mathbf{x} \,\mathrm{d}^3 \mathbf{v},$$

#### is *independent* of the assumed potential $\Phi$ .

We may average the DF over angles in an assumed  $\Phi$ :

$$ar{f}(\mathbf{J}) = rac{1}{(2\pi)^3} \int f(\mathbf{J}, heta) \mathrm{d}^3 heta.$$

Note that:

- $\bar{f} = f$  if assumed  $\Phi$  is true potential (because  $\theta$  flat)
- $S[\bar{f}] \ge S[f]$  because S[f] indep of  $\Phi$ .
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## But does it work?

#### Application to Gaia challenge problems

Apply to (**x**, **v**) from Matt W/Jorge P's spherical models in spherical\_df.tar.gz.

I assume a parametrised mass density of the form

$$\rho(\mathbf{r}) = \rho_0 \mathbf{r}^{-\gamma} (\mathbf{1} + \mathbf{r})^{\gamma - 3}.$$

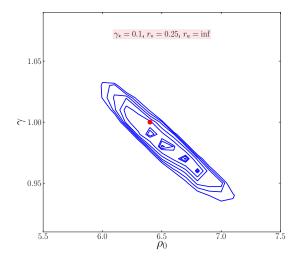
Matt's  $(\mathbf{x}, \mathbf{v})$  "data" come from models having

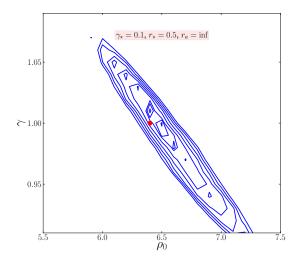
• 
$$\rho_0 = 6.4$$
 and  $\gamma = 1$ , but

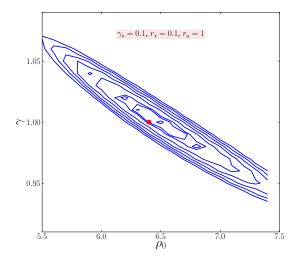
• a range of different stellar DFs (10<sup>4</sup> stars each).

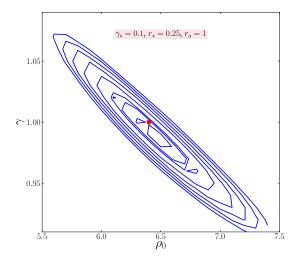
For each dataset *D* I estimate  $p(D|\rho_0, \gamma)$  using VB method.

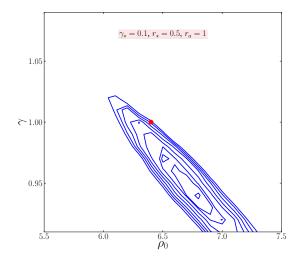
(NB: Following plots updated during the meeting: still not perfect, but VB is an *approximate* method...)

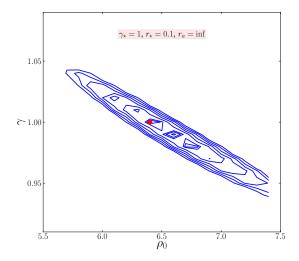


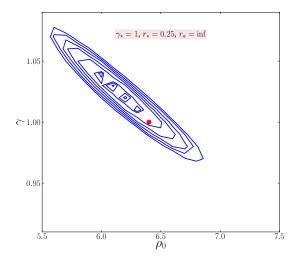


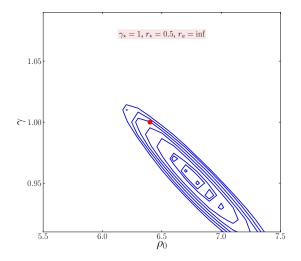


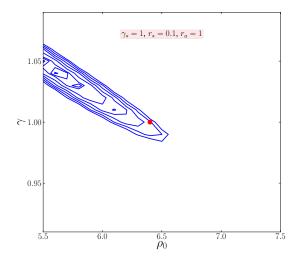


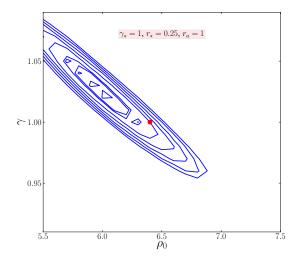


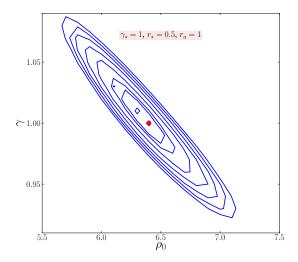












## Summary

DPM is sensible<sup>1</sup> way of representing prior on DFs:

- Models the DF "nonparametrically" (in true sense)
  - could view as unholy marriage of ∞-resolution Schwarzschild method in which δ-fn orbits replaced by parametrised DFs
  - (but that might not be helpful)
- Gives broadly correct results for  $p(D|\rho_0, \gamma)$ 
  - exact calculation possible only for  $N \leq 10$
  - approximate VB fast, but still room for improvement
  - beguiling interpretation in terms of streams.
- Straightforward in principle to include in likelihood
  - measurement errors (e.g., missing radial vels)
  - survey selection functions.
- Finding an effective practical method to do this is challenging.

<sup>&</sup>lt;sup>1</sup>Worth thinking about: do we *really* want to marginalise?